



Investigation of Shielding Parameters for Barite and Boron Carbide Polymer Composite

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Abstract

In this study, the shielding parameters for Glycidyl methacrylate with different ratio of Barite and Boron Carbide Polymer composites have been calculated at the photon energy range of 1 keV-1 GeV by using XCom program. The obtained data will be used to calculate the effective atomic number (Z_{eff}) and effective electron density (N_{eff}) for the same range of energy. The macroscopic fast neutron removal cross-sections have also been calculated. The fast neutron removal cross-sections on chemical composition of the selected polymer composites have been discussed. Also, the dependence of characteristics gamma-ray shielding parameters for the composite on incident photon energy has been studied. Barite shows up as good attenuating material, while Boron samples are relatively weak gamma-ray attenuators. The effect of the compound composition appears obviously in this study. As well as, the obtained results through this study can be utilized to comprehend the shielding effectiveness of this composite.

Keywords

Barite; Boron carbide; Polymer composite; Effective atomic number; Effective electron density; Fast neutron removal cross-sections

Introduction

Due to the increasing nuclear applications in different fields makes nuclear radiation to play an important role in our daily life. Various beneficial applications of different types of radiations are in the field of medicine, industry, agriculture, and research. Neutron and gamma-ray sources are not only dangerous to human being but also to sensitive laboratory equipment's. Radiation shielding is necessary to protect penetrative neutrons and gamma-rays. The intensity of penetrating radiation such as neutrons and gamma-rays is varied by three factors: time, distance and shielding. The most effective method for attenuation of radiation is shielding [1]. Filler-reinforced polymers have gained increasing attention from X-ray technologies in radiation shielding since polymers have great potential in many important applications by virtue of their unique properties, such as low density, the ability to form intricate shapes, optical transparency, low manufacturing cost, and toughness. One of the filler-reinforced polymers commonly used for radiation shielding is lead acrylic. Moreover, some researchers have also tried to synthesize nano-sized

filler-reinforced polymers for radiation shielding by virtue of the size effect in X-ray attenuation [2-4].

The mass attenuation coefficient, μ , is a measure of the average number of interactions between incident photons and matter that occur in a given mass-per-unit area thickness of the material encountered [5]. Accurate values of the mass attenuation coefficient, μ , are required in the nuclear diagnostics, radiation protection, nuclear medicine, radiation biophysics, radiation dosimetry to obtain essential data. The mass attenuation coefficients are also used to determine photon penetration and energy deposition in the biological and other materials [6-8]. The mass attenuation coefficients can be obtained from different database [9] and XCom program [10]. Because of photon interaction cross-section for elements of the composite material, a single atomic number is a characteristic of element will not describe the atomic number of composite material in all energy ranges. So a new number for composite materials is called to be effective atomic number, suggested by [9] and it varies with energy. Determination of the effective atomic number, Z_{eff} and the effective electron number per unit mass, N_{eff} which is another important parameter for composite materials is very important in the fields of many technological applications and nuclear medicine for the calculation of dose in radiation therapy and medical imaging [6,11].

Neutron penetration in shielding is characterized by several parameters such as the effective removal cross-sections, the macroscopic thermal neutron cross-section. Effective removal cross-section, Σ_r (cm²/g) is a measure of the probability that fast or fission energy neutron undergoes a first collision, which removes it from the group of the penetrative un-collided neutrons [12,13]. It is approximately constant for neutron energies between 2 and 12 MeV. The concept of Σ_r is used as long as the shielding material under investigation contains some scattering nuclides. However, when there are no scattering nuclides around the beam another quantity, the total mass neutron cross-section (cm²/g), is used. Observed values of the Σ_r are roughly 2/3 of Σ_T for neutrons of energies in the range from 6 to 8 MeV [14-17].

In this work, the total mass attenuation coefficient, μ_m has been determined theoretically by using XCom code at energies from 1 keV to 1 GeV. Since the effective atomic number, Z_{eff} and the effective electron number per unit mass, N_{eff} are widespread use; this study is focusing on calculating the effective atomic number and effective electron density of polymer composites with different barite and boron carbide percentages. Also, the effective removal cross-section has been calculated theoretically by using the values of removal cross-section of elements which constitute the polymer composition. The elemental composition for the Glycidyl methacrylate with different ratio of Barite and Boron Carbide Composite are displayed in Table 1.

Methodology of Calculations

Gamma-ray and neutron radiation shielding parameters for Barite and Boron Carbide Polymer composites were theoretically calculated as following:

Gamma-ray calculations

When a gamma-ray passes through a matter, its intensity

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Table 1: Abbreviation and elemental composition and density for glycidyl methacrylate with different ratio of barite and boron carbide composite.

Element	Fraction by weight																															
	Barite 0% (P1)	Barite 5% (PB2)	Barite 10% (PB3)	Barite 15% (PB4)	Barite 20% (PB5)	Barite 25% (PB6)	Barite 30% (PB7)																									
	!	J	F	P	!	J	F	P	!	J	F	P	!	J	F	P	!	J	F	P	!	J	F	P	!	J	F	P	!	J	F	P
H																																
C																																
O																																
S																																
Ba																																
Element	Boron carbide 0% (P1)	Boron carbide 5% (PBC2)	Boron carbide 10% (PBC3)	Boron carbide 15% (PBC4)	Boron carbide 20% (PBC5)	Boron carbide 25% (PBC6)	Boron carbide 30% (PBC7)																									
	!	J	F	P	!	J	F	P	!	J	F	P	!	J	F	P	!	J	F	P	!	J	F	P								
H																																
B																																
C																																
O																																
Element	Barites 0%	Barite 2.5%	Barite 5%	Barite 7.5%	Barite 10%	Barite 12.5%	Barite 15%																									
	Boron carbide 0% (P1)	Boron carbide 2.5% (PBBC2)	Boron carbide 5% (PBBC3)	Boron carbide 7.5% (PBBC4)	Boron carbide 10% (PBBC5)	Boron carbide 12.5% (PBBC6)	Boron carbide 15% (PBBC7)																									
	!	J	F	P	!	J	F	P	!	J	F	P	!	J	F	P	!	J	F	P												
H																																
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progressively reduces as a consequence of interactions between photons and atoms. The mass attenuation coefficient, (μ/ρ) , is a measure of the average number of interactions that occur between photons and matter.

Consider first a chemical compound. The formulas will later be generalized to mixtures as well. The total photon interaction cross-section, σ_m , per molecule can be written as follows:

$$\sigma_m = \sum_i n_i \sigma_i \tag{1}$$

Where n_i is the number of atoms of the constituent element, and σ_i is total photon interaction cross-section per atom of element i . The total number, n , of atoms in the molecule is

$$n = \sum_i n_i \tag{2}$$

Suppose that the cross-section per molecule can be written in terms of an effective (average) cross-section per atom and an effective cross-section per electron as follows:

$$\sigma_m = n \sigma_a = n Z_{eff} \sigma_e \tag{3}$$

From the first equality of Eq. (3) the effective cross-section per atom is given by:

$$\sigma_a = \frac{1}{n} \sum_i n_i \sigma_i \tag{4}$$

In the same way, the effective cross-section per electron is given by:

$$\sigma_e = \frac{1}{n} \sum_i n_i \frac{\sigma_i}{Z_i} \tag{5}$$

It follows from the last equality of Eq. (3) that the effective atomic number can be written as:

$$Z_{eff} = \frac{\sigma_a}{\sigma_e} \tag{6}$$

By inserting the expressions (4) and (5) for σ_a and σ_e :

$$Z_{eff} = \frac{\sum_i n_i \sigma_i}{\sum_i n_i \frac{\sigma_i}{Z_i}} \tag{7}$$

For a chemical compound one has

$$f_i = \frac{n_i}{\sum_j n_j} = \frac{n_i}{n} \tag{8}$$

Where $\sum_i f_i = 1$ Rewriting Eq. (7) in terms of molar fractions gives the more general expression

$$Z_{eff} = \frac{\sum_i f_i \sigma_i}{\sum_i f_i \frac{\sigma_i}{Z_i}} \tag{9}$$

Eq. (9) is then the basic relation for calculating the effective atomic number for all types of materials, compounds as well as mixtures.

The atomic cross-section, σ_i , of the i th constituent element, is related to the corresponding mass attenuation coefficient, $(\mu/\rho)_i$, through the relation:

$$\sigma_i = \frac{A_i}{N_A} \left(\frac{\mu}{\rho} \right)_i \tag{10}$$

Where A_i is the atomic mass and N_A is the Avogadro constant. Useful expressions in terms of the mass attenuation coefficient can be obtained by inserting expression (10) for the equations of the previous section.

Consider first a chemical compound. Inserting the expression (10) for σ_i in Eq. (7) gives the following relation for Z :

$$Z_{eff} = \frac{\sum_i f_i A_i \left(\frac{\mu}{\rho} \right)_i}{\sum_i f_i \frac{A_i}{Z_i} \left(\frac{\mu}{\rho} \right)_i} \tag{11}$$

Eq. (11) can be used for calculating the effective atomic number Z_{eff} . A material sample having as components: simple elements both compounds and mixtures.

The effective atomic number, Z_{eff} is closely related to the electron density, N_e , which is expressed in a number of electrons per unit mass. For a chemical element, the electron density is given by $N_e = \frac{N_A Z}{A}$. This expression can be generalized to a compound, and one has

$$N_e = N_A \frac{\sum n_i Z_i}{\sum n_i A_i}$$

$$N_A = \frac{Z_{eff}}{\langle A \rangle} \quad (12)$$

Where $\langle A \rangle$ is the average atomic mass of the compound. It can be shown that the electron density also is given by

$$N_{eff} = \frac{\left(\frac{\mu}{\rho} \right)}{\sigma_e} \quad (13)$$

Where μ is the total mass attenuation coefficient of the compound, and σ_e is the electronic cross-section given by Eq. (5). In the case of Compton scattering, the Klein-Nishina cross-section can be used for σ_e .

Fast neutron calculations

The interaction of neutrons with the atoms described by the total microscopic cross-section, σ_t , expresses the probability that a neutron of given energy interacts with the atoms of the traversed material and it is defined as the sum of the microscopic cross-section of scattering (σ_s) and the microscopic cross-section absorption (σ_a)

$$\sigma_t = \sigma_s + \sigma_a \quad (14)$$

The neutrons attenuation during their passage through a material medium depends not only on the microscopic cross-section but also on the number of nuclei within this environment. The physical quantity bounding these two parameters is called total macroscopic cross-section denoted Σ_t and defined by [18,19]:

$$\Sigma_t = \frac{\rho N_A \sigma_t}{A} \quad (15)$$

Where ρ is the density (g/cm^3), N_A is Avogadro's number and A is the atomic mass. In the same way as a photons beam, when the parallel beam of monoenergetic neutrons passes through a material medium, it will be attenuated due to absorption and scattering. The attenuation of neutrons in matter follows the following law:

$$I = I_0 e^{-\Sigma_t x} \quad (16)$$

where I_0 and I are the incidents and transmitted intensities, x (cm) is the thickness of the material medium and Σ_t represents the total macroscopic cross-section.

So the case of fast neutron attenuation is described by another parameter called the removal cross-section, denoted Σ_R (cm^{-1}) and is different from the total macroscopic cross-section but it has a fraction of it. The removal cross-section presents the probability that a fast or fission-energy neutron undergoes a first collision, which removes it from the group of penetrating un-collided neutrons.

For energies between 2 and 12 MeV, the effective removal cross-section will be almost constant when the traversed medium contains a large amount of hydrogen $\Sigma_R = \sum \sigma_{R,i}$ and when materials contain a small fraction of hydrogen $\Sigma_R = \frac{2}{3} \sum \sigma_{R,i}$ for energy between 6 and 8 (0.1 > E < 8 MeV). It is clear that Σ_R increases with increasing of energy

and compounds, its removal cross-section is given by the following formula [14-17]:

$$\Sigma_R = \sum_i W_i \left(\frac{\Sigma_{R,i}}{\bar{n}_i} \right) \quad (17)$$

Where W_i , ρ_i and $\left(\frac{\Sigma_{R,i}}{\bar{n}_i} \right)_i$ are respectively the partial density (g/cm^3), density and mass removal cross-section of the i th constituent. The partial density of the i th constituent is given by:

$$W_i = (t_i) \rho_s \quad (18)$$

Where (t_i) is the weight fraction of the i th constituent and ρ_s is the sample density.

In this study, the effective removal cross-section (Σ_R) of fast neutrons has been calculated for polymer composites with different ratios of barite and boron carbide by using Eqn. 17. The values of the mass removal cross-sections of the elements that constitute of these materials are taken from [20,21].

Results and Discussion

In this study, the basic radiation parameters of polymer composites with different ratios of Barite and Boron Carbide were studied by the direct method over a wide photon energy range from 1 keV to 1 GeV using XCom program. These parameters are the μ/ρ , Z_{eff} . As well as the effective removal cross-section (Σ_R) of fast neutrons has been calculated for the investigated Samples.

The Z_{eff} and N_e values of Polymer composites (equal proportions of barite and boron carbide) are given in Tables 2 and 3 only for total photon interaction. The tables present Z_{eff} and N_e values in the energy range from 1 keV to 20 MeV, however, the values for the energy range from 1 keV to 1 GeV are displayed in graphical results. The elemental composition of materials used in this work, its fractions by weight and calculated $\left(\frac{\Sigma_{R,i}}{\bar{n}_i} \right)$ values are listed in Table 4.

Mass attenuation coefficients

The dependence of mass attenuation coefficient μ/ρ of polymer composites to photon energy is shown in Figures 1-3. While in Figure 4 only results from (PB5) sample on the variations of the μ/ρ for total and partial photon interactions were shown. It is seen from the figure that different dominant interaction processes on different energy regions. There are three energy ranges relative to the partial processes: photoelectric absorption at low energies, incoherent (Compton) scattering at intermediate energies and pair production at high energies. In Figure 4, the curve of Compton scattering is split up because it dependent on probability of photon interaction by incoherent scattering for each energy.

Effective atomic number, Z_{eff}

The value of Z_{eff} for a composite material is a very useful parameter in some applications such as physical, technological and engineering. This value can provide an estimation of the chemical composition of the material and it can be also utilized in the computation of absorbed dose in radiation therapy etc. In this work, Z_{eff} values for Polymer composites (equal proportions of barite and boron carbide) were calculated by using Eqn. (11) and given in Table 2.

The variation of Z_{eff} of photon energy for total interaction processes in all-polymer composites is shown in Figures 5-7. These figures show that the minimum value of Z_{eff} is found around intermediate energies (0.1 > E < 8 MeV). It is clear that Z_{eff} increases with increasing of energy

Effective removal cross-sections of fast neutrons

The results of calculated Σ_{Rf} for the investigated composites are given in Table 4. It can be seen from these tables that the obtained values for Σ_{Rf} are ranged from 6.96E-02, for (PB7) sample, to 8.58E-02, for (P1) sample. The results obtained were compared with the special values of three types of concrete [22] as shown in Table 4, which showed a relative increase in the values investigated

samples. This means that the investigated samples are more effective as fast neutron attenuators than the common materials used in that field. Figure 12 has shown the calculated effective removal cross-sections of fast neutrons for polymer composites as a function of fraction by weight for each sample. It was found that Σ_{Rf} depends on the elemental composition and it is clear that the most important parameter affecting neutron attenuation is the amount of hydrogen in the sample. This can be attributed to the fact that the hydrogen has a higher value $\sigma_{Rf}(E)$.

As shown from Figure 12, the higher concentration of hydrogen in the chemical composition of the polymer without any addition in comparison to that of the other samples with different ratios of barite and boron carbide explain the greater value of Σ_{Rf} for the polymer in comparison to values of Σ_{Rf} for other samples. While the low value of Σ_{Rf} observed for (PB7) sample is mainly due to the low concentration of hydrogen in the material. Hence, it is concluded that the concentration of hydrogen in the polymer is the most important parameter that affects the attenuation of fast neutrons in such materials.

At first sight, we can conclude from these results that the (P1) sample is the most effective material for neutron shielding, which is not true. For an effective attenuation of fast neutrons in the neutron shielding material, both moderator and absorber should exist in the

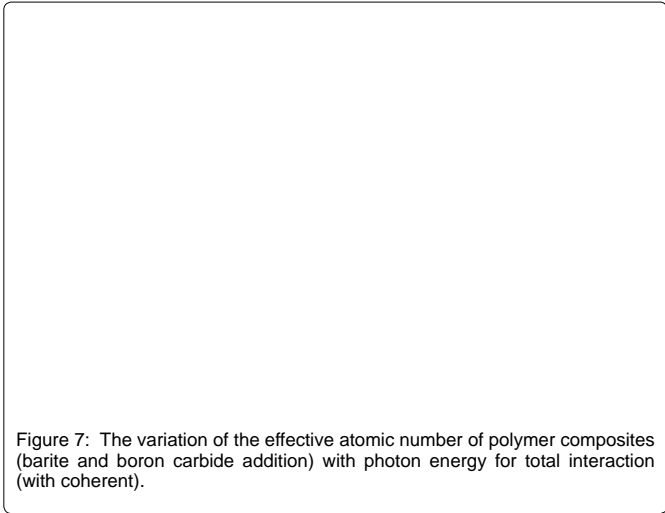


Figure 7: The variation of the effective atomic number of polymer composites (barite and boron carbide addition) with photon energy for total interaction (with coherent).

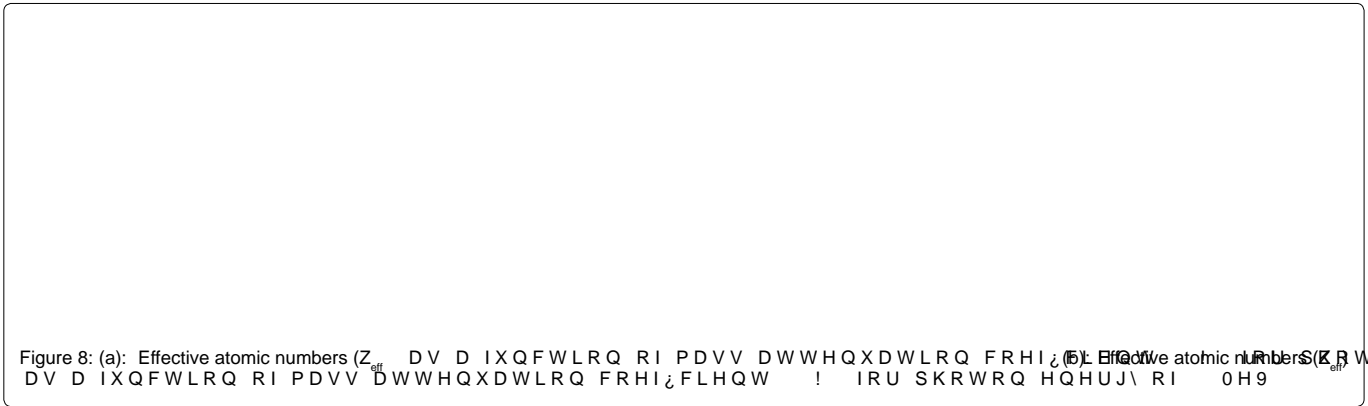


Figure 8: (a): Effective atomic numbers (Z_{eff}) and (b): Effective atomic numbers (Z_{eff}) for different samples.

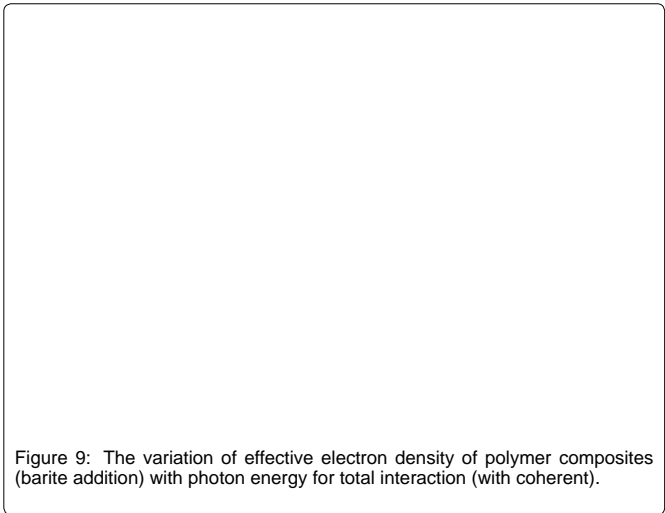


Figure 9: The variation of effective electron density of polymer composites (barite addition) with photon energy for total interaction (with coherent).

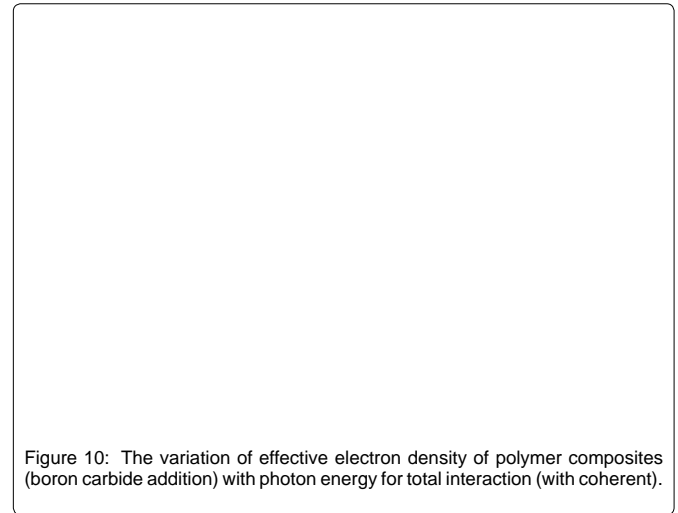


Figure 10: The variation of effective electron density of polymer composites (boron carbide addition) with photon energy for total interaction (with coherent).

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