# From Configurations to Branched Configurations and Beyond 

Jean-Pierre Magnot*


#### Abstract

We propose here a geometric and topological setting for the study of branching effects arising in various fields of research, e.g. in statistical mechanics and turbulence theory. We describe various aspects that appear key points to us, and finish with a limit of such a construction which stand in the dynamics on probability spaces where it seems that branching effects can be fully studied without any adaptation of the framework.


## Keywords

Finite and infinite configuration spaces; Branched dynamics

## Introduction

Finite and infinite configuration spaces are rather old topics, see e.g. [ 1,2 ] and the references cited therein, that had many applications in various settings in mathematical physics and representation theory. More recently, several papers, including [3-6] showed how these topics could be applied in various disciplines: ecology, financial markets, and so on. This large spectrum of applications principally comes from the simplicity of the model: considering a state space N , finite or infinite configurations are finite or countable sets of values in N : This is why we begin with giving a short description of this setting, and describe a differentiable structure that can fit with easy problems of dynamics. This structure, which can be seen either as a Frölicher structure or as a diffeological one, is carefully described and the links between these two frameworks are summarized in the appendix. We also give a result that seems forgotten in the past literature: the infinite configuration space used in e.g. [7] is an infinite dimensional manifold.

But the main goal of this paper is to include one dimensional turbulence effects (in particular period doubling, see e.g. [8]) in the dynamics described by finite or infinite configuration spaces. For this, we change the metric into the Hausdorff metric. This enables to "glue together" two configurations into another one and to describe shocks. Therefore, the dynamics on this modified configuration space are described by multivalued paths that are particular cases of graphs on $N$; which explains the terminology: "branched configuration".

We finish with the description of the configurations used e.g.

[^0]in optimal transport, the space of probability measures, and shows how they can also furnish configurations for uncompressible fluids. In these settings, branching effects are well-known and sometimes obvious, and we do not need any adaptation of the framework to obtain a full description of them. Therefore, what we call branched configurations appears as an intermediate (an we hope useful) step between dynamics of e.g. a N -body problem and e.g. wave dynamics.

## Dirac configuration spaces on a locally compact manifold

Let us describe step by step a way to build infinite configurations, as they are built in the mathematical literature. We explain each step with the configurations already defined in e.g. Albeverio et al. [7] and Fadell et al. [1], the generalization will be discussed later in this paper.

A set of $I$-configurations is a set of objects that are modelizations of physical quantities. For example, in the settings [1-6], the physical quantity modeled is the position of one particle. The whole world is modeled as a locally compact manifold $N$, and the set of 1 -configurations is itself $N$, or equivalently the set of all Dirac measures on $N$.

Let $I$ be a set of indexes. $I$ can be countable or uncountable. We define the indexed (or the ordered if $I$ isequipped with its total order) configuration spaces. For this, we need to define a symmetric binary relation relation $u$ on $\Gamma^{1}$, that expresses the compatibility of two physical quantities. We assume also that $u$ has the following property:

$$
\forall(u, v) \in\left(\Gamma^{1}\right)^{2}, u U v \Rightarrow u \neq v
$$

In the settings Albeverio et al. [7] and Finkelshtein et al. [1], two particles cannot have the same position. Then, for $x, y \in N[2]$,

$$
x u y \Leftrightarrow x \neq y
$$

With these restrictions, we can define the indexed or ordered configuration spaces:

$$
\begin{aligned}
& O \Gamma^{n}=\left\{\left(u_{1}, \ldots, u_{n}\right) \in\left(\Gamma^{1}\right)^{n} \text { such that, } i f i \neq j ; u_{1} u u_{j}\right\} \\
& O \Gamma=\coprod_{n \in \mathbb{N}^{*}} O \Gamma^{n} \\
& O \Gamma^{I}\left\{\left(u_{n}\right)_{n \in I} \in\left(\Gamma^{I}\right)^{I} \text { such that, ifi} \neq j ; u_{1} u u_{j}\right\}
\end{aligned}
$$

The general configuration spaces are not ordered. Let $\Sigma_{n}$ be the group of bijections on $\mathbb{N}_{n}$, and $\Sigma_{I}$ be the set of bijections on $I$. We can define two actions:

$$
\begin{aligned}
& \sum_{n} \times O \Gamma^{n} \rightarrow O \Gamma^{n} \\
& \left(\sigma,\left(u_{1}, \ldots, u_{n}\right)\right) \rightarrow\left(u_{\sigma}(1), \ldots, u_{\sigma}(\mathrm{n})\right)
\end{aligned}
$$

and its infinite analog:

$$
\begin{aligned}
& \sum_{I} \times O \Gamma^{I} \rightarrow O \Gamma^{I} \\
& \left(\sigma,\left(u_{n}\right)_{n \in I}\right) \rightarrow\left(u_{\sigma(n)}\right)_{n \in I}
\end{aligned}
$$

where $\Sigma_{I}$ is a subgroup of the group of bijections of $I$ : In the sequel, $I$ is countable with discrete topology, which avoids topological problems on $\Sigma_{I}$ as in more complex examples. Then, we define general configuration spaces:

$$
\Gamma^{n}=O \Gamma^{n} / \sum_{n}
$$

$$
\begin{aligned}
& \Gamma=\coprod_{n \in \mathbb{N}^{*}} \Gamma^{n} \\
& \Gamma^{I}=O \Gamma^{I} / \Sigma_{I}
\end{aligned}
$$

## Configuration spaces on more general settings

In the machinery of the last section, the properties of the base manifold $N$ are not used in the definition of the space $\Gamma^{1}$. This is why the starting point can be $\Gamma^{1}$ instead of $N$, and we can give it the most general differentiable structure. Let us first consider the most general case [9]:

Proposition 1.1: If $\Gamma^{1}$ is a diffeological space, then $\Gamma^{n}, \Gamma$ and $\Gamma^{I}$ are diffeological spaces

Proof: $\left(\Gamma^{1}\right)^{n}\left(\right.$ resp. $\left.\left(\Gamma^{1}\right)^{\mathbb{N}}\right)$ is a diffeological space according to Proposition 4.8 (resp. Proposition 4.10), so that, $O \Gamma^{n}\left(\right.$ resp. $\left.O \Gamma^{1}\right)$ is a diffeological space as a subset of $\left(\Gamma^{1}\right)^{n}$ (resp. $\left.\left(\Gamma^{1}\right)^{N}\right)$. Thus, $\Gamma^{n}=O^{n} /$ $\Sigma_{n}\left(\right.$ resp. $\left.\Gamma^{1}=O^{n} / \Sigma\right)$ has the quotient diffeology by Proposition 4.12, which ends the proof.

Let us now turn to the cases where $\Gamma^{I}$ has a stronger structure. We already know that $\Gamma^{n}$ is a manifold if $\Gamma^{1}$ is a manifold.

Proposition 1.2: If $\Gamma^{1}$ is a Frölicher space, then $\Gamma^{n}$ (and hence $\Gamma$ ) is a Frölicher space.

Proof: Adapting the last proof, using Proposition 4.9 instead of Proposition 4.8, we get that $O \Gamma^{n}$ is a Frolicher space if $\Gamma^{1}$ is a Frolicher space. Let us now build a generating set of functions for the Frolicher structure on $O \Gamma^{n}$. Let $f: O \Gamma^{n} \rightarrow \mathbb{R}$ be a smooth map. We define the symmetrization of $f$ :

$$
\tilde{f}:\left(u_{1}, \ldots, u_{n}\right) \in O \Gamma^{n} \mapsto \tilde{f}\left(u_{1}, \ldots, u_{n}\right)=\frac{1}{n!} \sum_{\sigma \in \Sigma_{n}} f\left(u_{\sigma(1)}, \ldots, u_{\sigma(n)}\right)
$$

The set of functions $\{\tilde{f}\}$ generate the contours on $O \Gamma^{n}$, and then, passing to the quotient, generate the contours on $\Gamma^{n}$

Setting a Frölicher structure on indexed configurations is only a straightforward consequence of Proposition 4.10:

Proposition 1.3: If $\Gamma^{1}$ is a Frölicher space, then $O \Gamma^{1}$ is a Frölicher space.
But the problem of a Frölicher structure on $\Gamma^{1}$ a little bit more complicated; let us explain why and give step by step the construction of the Frölicher structure. We first notice that the contours of the Frölicher push forward naturally by the quotient map $O \Gamma^{1} \rightarrow \Gamma^{1}$. But, if one wants to describe a generating set of functions, by Proposition 4.9, one has to consider all combinations of a finite number of smooth functions $\Gamma^{1} \rightarrow \mathbb{R}$ : This generating set does not contain any $\Sigma_{I}$-invariant function, except constant functions. This is why the approach used in the proof of the Proposition 1.2 cannot be applied here. For this, one has to consider the set

$$
F_{e q}=\left\{f: O \Gamma^{I} \rightarrow \mathbb{R} \in F \text { such that } \mathrm{f} \text { is } \Sigma^{I} \text {-invariant }\right\}
$$

of equivariant functions on $O \Gamma^{\mathrm{I}}$ : This discusssion can become very quickly naïve and we prefer to leave this question to more applied works in order to fit with known examples instead of dealing with too abstract considerations.

## Topological Configuration Spaces

In this section, we present examples of 1-configurations and their associated configuration spaces. Manifolds will replace the Dirac measures used in Albeverio et al. [7]. In the sequel, $N$ is a Riemannian
smooth locally compact manifold. The 1-configurations considered keep their topological properties, as in the model of elastodymanics (see e.g. Hughes et al. [10]) or in various quantum field theories. Notice also that we do not give compatibility conditions between two 1-configurations: we would like to give the more appropriate conditions in order to fit with the applied models, this is why we leave this point to more specialized works.

## Topological 1-configurations

We follow here, for example, Hughes T et al. [10].
Definition 2.1: Let $M$ be a smooth compact manifold and $N$ an arbitrary manifold. We set

$$
\Gamma_{e}^{I}(M, N)=C^{\infty}(M, N)
$$

One can also only consider embeddings, and set:

$$
\Gamma_{m}^{I}(M, N)=\operatorname{Emb}(M, N)
$$

where $\operatorname{dimM}<\operatorname{dim} N$, and Emb is the set of embeddings. The things run as in the first case, since $\operatorname{Emb}(M ; N) \subset C^{\infty}(M ; N)$ is an open subset of $C^{\circ}(M ; N)$ :

## Examples of topological configurations

Links. Let $\Gamma^{1}=\operatorname{Emb}\left(S^{1} ; N\right)$ Here, we fix the uncompatibility relation as

$$
\gamma u \gamma^{\prime} \Leftrightarrow \gamma\left(\mathrm{S}^{1}\right) \cap \gamma^{\prime}\left(\mathrm{S}^{1}\right) \neq \varnothing
$$

Then, $\Gamma_{\text {linkk }}^{n}$ is the space of $n \rightarrow$ links of class $C^{k}$, which is a Frechet manifold.

Triangulations: Consider the n -simplex

$$
\Delta_{n}=\left\{\left(t, \ldots, t_{n}\right) \in \mathbb{R} \quad \mid \sum t_{i}=1\right\}
$$

If $N$ is a n -dimensional manifold, a (finite) triangulation $\sigma$ of $N$ is such that:

1) $\sigma \in\left(\Gamma_{m}^{1}\left(\Delta_{n}, N\right)\right)^{|\sigma|}$
(2) Let $\left(T, T^{\prime}\right) \in \sigma^{2}$ such that $T \neq T^{\prime}$ then $\operatorname{Im}(T) \cap \operatorname{Im}\left(T^{\prime}\right)$ is a simplex or a collection of simplexes of each border $\operatorname{Im}(\partial T)$ and $\operatorname{Im}\left(\partial T^{\prime}\right)$ :
(3) $\bigcup_{T \in \sigma} \operatorname{Im} T=N$

We get by condition 2 a compatibility condition $u$; for which we can build $O \Gamma\left(\Delta_{n}, N\right), \Gamma\left(\Delta_{n}, N\right)$; $O \Gamma^{\infty}\left(\Delta_{n^{\prime}}, N\right)$ and $\Gamma^{\infty}\left(\Delta_{n}, N\right)$; If $N$ is compact, the set of triagulations of $N$ is a subset of $\Gamma\left(\Delta_{\mathrm{p},} N\right)$ : If N is non compact and locally compact, the set of triangulations of $N$ is a subset of $\Gamma^{\circ}\left(\Delta_{n^{\prime}} N\right)$ :

More generally, for $p \leq n$; one can build $O \Gamma\left(\Delta_{\mathrm{p}}, N\right), \Gamma\left(\Delta_{\mathrm{p}}, N\right), O \Gamma^{\infty}\left(\Delta_{p}\right.$ $N)$ and $\Gamma^{\circ}\left(\Delta_{p} N\right)$ : This example will be discussed in the section 3.

Strings and membranes: A string is a smooth surface $\Sigma$ possibly with boundary, embedded in $\mathbb{R}^{26}$ : A membrane is a manifold $M$ of higher dimension embedded in some $\mathbb{R}^{k}$ : We recover here some spaces of the type $\Gamma_{m}^{1}$, which will be also discussed in section 3.

## Branched Configuration Spaces

## Dirac branched configurations

As we can see in section 1.1, finite configurations $\Gamma$ Tare made of a countable disjoint union. We now fix a metric $d$ on $N$ : The idea of
branched configurations is to glue together the components $\Gamma^{n}$ on the generalized diagonal, namely, we define the following distance on $\Gamma$ :

Definition 3.1: Let $(u ; v) \in \Gamma^{2}$ :

$$
d_{\Gamma}(u, v)=\sup _{\left(x, x^{\prime}\right) \in u \times v}\left\{d(\mathrm{x}, v), d\left(\mathrm{x}^{\prime}, \mathrm{u}\right)\right\}
$$

Proposition 3.2: $d_{\Gamma}$ is a metric on $\Gamma$ :
Proof: We remark that $d_{\Gamma}$ is the Hausdorff distance restricted to $\Gamma$ :
The following proposition traduces the change of topology of $\Gamma$ into $B \Gamma$ by cutand paste property:

Proposition 3.3: $\forall n \in \mathbb{N}^{*}, B \Gamma^{\mathrm{n}+1}=\Gamma^{\mathrm{n}+1} \coprod_{B \Gamma^{n+1}} B \Gamma^{\mathrm{n}}$ where the identification is made along the trace on $\Gamma^{\mathrm{n}+1}$ of the $d_{\Gamma}-$ neighborhoods of $B \Gamma^{\mathrm{n}}$ in $N^{n+1} \supset O \Gamma^{\mathrm{n}+1}$ :

We remark that we can also define a Frolicher structure on $B \Gamma$ with generating set of functions the set

$$
\left\{\left.u \in \varnothing \Gamma \mapsto \frac{1}{|u|} \sum_{x \in u} f(x) \right\rvert\, f \in C^{\infty}(N, \mathbb{R})\right\}
$$

This structure will be recovered later in this paper.

## Examples of topological branched configurations

The path space, branched paths and graphs: Let

$$
\Gamma_{e}^{1}\left([0 ; 1] ; N=C^{\infty}[0 ; 1] ; N\right)
$$

be the space of smooth paths on $N$ : A $\gamma$ path has a natural orientation, and has a beginning $\alpha(\gamma)$ and an end $\omega(\gamma)$ : We define a compatibility condition

$$
\gamma u \gamma^{\prime} \Leftrightarrow\left(\left(\operatorname{Im} \gamma \neq \operatorname{Im} \gamma^{\prime} \vee(\alpha(\gamma), \beta(\gamma)) \neq\left(\alpha\left(\gamma^{\prime}\right), \beta\left(\gamma^{\prime}\right)\right)\right)\right.
$$

and we remark that the set of piecewise smooth paths on $N$ is a subset of $O \Gamma_{e}([0 ; 1] ; N)$; saying that $\left(\gamma_{1}, \ldots . \gamma_{p}\right) \in o \Gamma_{e}^{p}([0 ; 1] ; N)$ is a piecewise smooth path if and only if

$$
\forall i \in \mathbb{N}_{p-1}, \omega\left(\gamma_{i}\right)=\left(\gamma_{i}+1\right)
$$

This relation, stated from the natural definition of the composition * of the groupoid of paths, is not unique and can be generalized.

Definition 3.4: Let $\left(\left(\gamma_{1} ; \gamma_{2}\right) \in o \Gamma_{e}^{2}([0 ; 1] ; N)\right.$ and let $\gamma_{3} \in \mathrm{C}^{\infty}([0 ; 1]$; $N)$ : We define the equivalence relation $\sim$ by

$$
\left(\gamma_{1}, \gamma_{2}\right) \sim \gamma_{3} \Leftrightarrow \gamma_{3}=\gamma_{2}^{*} \gamma_{1}
$$

The maps $\alpha: \gamma \mapsto \alpha(\gamma)$ and $\omega: \gamma \mapsto \omega(\gamma)$ extends to \set theorical" maps $o \Gamma_{e}^{k}([0 ; 1] ; N) \rightarrow N^{\mathrm{k}}$ and $\Gamma_{\mathrm{e}}([0 ; 1] ; N) \rightarrow(N)$ : The following is now natural:

Definition 3.5: A branched path is an element $\gamma$ of $O \Gamma_{e}\left(O \Gamma_{e}([0 ;\right.$ $1] ; N) / \sim)$ such that, if $\gamma O \Gamma^{\mathrm{k}}\left(O \Gamma_{\mathrm{e}}([0 ; 1] ; N) / \sim\right)$,

$$
\forall i \in \mathbb{N}_{k-1}, \omega\left(\gamma_{i}\right)=\left(\gamma_{i}+1\right) \text { in } \Gamma(\mathrm{N})
$$

Example: Let us consider the following paths $[0 ; 1] \rightarrow \mathbb{R}^{2}$ :
$\gamma_{1}(t)=(t-2 ; 0)$
$\gamma_{2}(t)=(\cos (\pi(1-t)) ; \sin (\pi t))$
$\gamma_{3}(t)=(\cos (\pi(1-t)) ;-\sin (\pi t))$
$\gamma_{4}(t)=(t+1 ; 0)$
Then,
$\omega\left(\gamma_{1}\right)=\alpha\left(\gamma_{2}\right)=\alpha\left(\gamma_{3}\right)$

And
$\omega\left(\gamma_{2}\right)=\omega\left(\gamma_{3}\right)=\left(\gamma_{4}\right)$
This shows that

$$
\left(\gamma_{1},\left(\gamma_{2}, \gamma_{3}\right), \gamma_{4}\right) \in O \Gamma^{3}\left(O \Gamma_{e}\left([0 ; 1] \rightarrow \mathbb{R}^{2}\right)\right.
$$

is a branched path of $\mathbb{R}^{2}$ :
Alternate approach to branched paths: branched sections of a fiber bundle.

Let $\pi: F \rightarrow M$ be a fiber bundle of typical fiber $F_{0}$ :
Here, $n \in \mathbb{N}^{*} U_{\infty}$ Let $\pi: F \rightarrow M$ be a fiber bundle over $M$ with typical fiber $F_{0}$ : Let

$$
\Gamma_{m}^{n}(\mathrm{~F})=\left\{\mathrm{u} \in \Gamma^{n}(F) \| \pi(u) \mid=1\right\}
$$

This is trivially a fiber bundle of basis $M$ with typical fiber $\Gamma^{n}\left(F_{0}\right)$ :
Definition 3.6: A non-section of $F$ is a section of $\Gamma_{m}^{n}(F)$ which cannot be decomposed into $n$ sections of $F$.

We define also $\Gamma_{\mathrm{M}}(F)=\amalg_{n \in \mathbb{N}^{*}} \Gamma_{m}^{n}(F)$; and also $\Gamma_{m}^{I}(F)$ the non sections based of $\Gamma^{I}\left(F_{0}\right)$ : We can define the same way $B \Gamma_{M}(F)$ using the branxhed configuration space instead of the configuration space, since the definitions from the set-theoric viewpoint are the same.
(Toy) Example: Let us consider the following example: $\mathrm{X}=\mathbb{R}^{3} \mathrm{X}$ \{up; down\}, and $\Gamma^{1}(\mathrm{X})=\mathbb{R}^{3}\{\{$ up; \{down\}; \{up; down\}\}; that models the position of an electron in the 3-dimensional space $\mathbb{R}^{3}$, associated to its spin. When the electron spin cannot be determined (i.e. out of the action of adequate electromagnetic fields), the picture proposed by Schrodinger is to consider that its spin is both up and down (this picture is also called the \Schrodinger cat" when we replace lup" and \down" by \dead" and \alive").

Let us now consider the Frolicher structure described on section 3.1. It is based on the natural diffeology carried by each $\Gamma^{n}\left(F_{0}\right)$ $\left(n \in \mathbb{N}^{*}\right)$ and by the set of paths $P_{1}^{\prime}$ that are paths $\gamma: \mathbb{R} \rightarrow \Gamma\left(F_{0}\right)$ such that $\exists(m, n) \in\left(\mathbb{N}^{*}\right)^{2}$
$\left.\gamma\right|_{-\times ; 0]}$ is a smooth path on $\Gamma^{m}(X)$
$\left.\gamma\right|_{\mathrm{j} 0 ;+\infty ; 0}$ is a smooth path on $\Gamma^{n}(X)$
Let $l \in \gamma(0)$. Then for any smooth map $f: F \rightarrow \mathbb{R}$
where $c$ are the trajectories going to $l$ in 0 equals to the sum of the infinite jet of

$$
\sum f \circ c_{-}
$$

where $c+$ are the trajectories coming from $l$ in $0^{+}$:

$$
\sum f \circ c_{+}
$$

Remark that the last condition comes from the smoothness required for each map $f \circ \gamma$ with $f \in C^{\infty}\left(F_{0} ; R\right)$ : This fits with the (fiberwise) frolicher structure of $B \Gamma\left(F_{0}\right)$ : Then, a finitely branched section of F is a smooth section of $B \Gamma_{\mathrm{M}}(F)$ : The first examples that come to our mind are the well-known branched processes, and we can wonder some deterministic analogues replacing stochastic processes by dynamical systems. Let us here sketch a toy example extracted from the theory of turbulence:

Example: equilibum of mayies population Assuming that Mayies live and die in the same portion of river, the population $p_{n+1}$ at the year $(n+1)$ is obtained from the population $p_{n}$ at the year n (after
normalization procedure) by the formula

$$
p_{n+1}=A_{p n}\left(1-p_{n}\right)
$$

where $A \in[0 ; 4]$ is a constant coming from the environmental data. For $A$ small enough, the fixed point of the so-called "logistic map" $\varphi_{A}(x)=A x(1-x)$ is stable, hence the population $p_{\mathrm{n}}$ tends to stabilize around this value. But when $A$ is increasing, the fixed point becomes unstable and $p_{\mathrm{n}}$ tends to stabilize around $2^{\mathrm{k}}$ multiple values which are the stable fixed points of the map $\phi^{2^{k}}$ obtained by composition rule.

Now, assume that we consider a river (or a lake), modelized by an interval (or an open subset of $\mathbb{R}^{2}$ ) that we denote by $U$, where the parameter $A$ is a smooth map $U \rightarrow[0 ; 4]$ : The parameter $A$ is a smooth map $U \rightarrow[0 ; 4]$ and the cardinality of configuration of equilibrum depends on the value of $A$ :

Non sections in higher dimensions: The example of a lake where mayflies live and die gives us a nice example of branched surface viewed as an element of $B \Gamma_{\mathbb{R}^{2}}\left(\mathbb{R}^{2} \times[0 ; 1]\right)$ : The same procedure can be implemented in gluing simplexes, or strings or membranes along their borders to get branched objects, but we prefer to postpone this problem to a work in progress where links with stochastic objects should be performed.
Measure-Like Configurations: An Example at the Borderline of Branched Configurations and Dynamics on Probability Spaces

Dynamics on probability space is a fast-growing subject and is shown to give rise to branched geodesics [11]. Following the same procedure as for branched topological configurations, we show her how a restricted space fits with particular goals. The goals described here are linked with image recognition for the configuration space ${ }^{h}$ above, and to uncompressible fluid dynamics when we equip $\Gamma_{h}^{1}$ with the diffeology $P_{0}$ of constant volume above. Let $C_{c}^{0}$ be the set of compactly supported $\mathbb{R}$-valued smooth maps on $N$. We define the relation of equivalence $R$ by:

$$
f R g \Leftrightarrow \sup p(f)=\sup p(g)
$$

Let us first give the definition of the set of 1-configurations:

$$
\Gamma_{h}^{1}(N)=C_{c}^{0}(N) / \mathrm{R}
$$

Definition 4.1: We set
Such a space is not a manifold, but we show that it carries a natural diffeology. $C_{c}^{0}(N)$ is a topological vector space, and hence carries a natural diffeology $P_{0}$. We define the following:

Definition 4.2: Let $P_{1} \subset P_{0}$ be the set of $P_{0}$-plots $p: O \rightarrow C_{c}^{0}(N)$ such that, for any open subset $A$ with compact closure $\overline{\mathrm{A}}$ of $N$, for any open subset $O^{\prime}$ of $O$ such that $\bar{O} \subset O$, if

$$
\mathrm{p}\left(O^{\prime}\right)(\bar{A} \backslash A)=0
$$

the map

$$
x \in O \mapsto \operatorname{Vol}(\sup p(p(x)) \cap A)
$$

is constant on $O^{\prime}$, where Vol is the Riemannian volume.
This technical condition ensures that the volume of any connected component of the support of $p(x)$ is constant. $P_{1}$ is obviously a diffeology on $C_{c}^{0}(N)$, and we can state

Proposition 4.3: Let $P=P_{1} / R .\left(\Gamma_{h}^{1}(N), \mathrm{p}\right)$ is a diffeological space.

The proof is a straightforward application of Definition 4.12. It seems difficult to give this space a structure of Frölicher space, or even a natural topology except the topology of vague convergence of measures, which is not the topology induced by the diffeology we have defined. As a consequence, we can only state that the welldefined configuration spaces $\Gamma$ and $\Gamma^{\Gamma}$ are diffeological spaces. The technical conditions of Definition 4.2 ensures that volume preserving is a consequence of smoothness with respect to $P_{1}$; and hence is particularily designed for (viscous) uncompressible fluid dynamics. A 1-confiuguration can have many connected components, and therefore branching effects are included in the definition of 1 -configurations. One could understand $n$-configurations as the presence of $n$ (non mixing) fluids.

## Appendix: Preliminaries on Differentiable Structures

The objects of the category of -finite or infinite- dimensional smooth manifolds is made of topological spaces $M$ equipped with a collection of charts called maximal atlas that enables one to make differentiable calculus. But there are some examples where a differential calculus is needed whereas no atlas can be defined. To circumvent this problem, several authors have independently developped some ways to define differentiation without defining charts. We use here three of them. The first one is due to Souriau [12], the second one is due to Sikorski, and the third one is a setting closer to the setting of differentiable manifolds is due to Frölicher (see e.g. Cherenack P et al. [13] for an introduction on these two last notions). In this section, we review some basics on these three notions.

## Souriau's DiffeologicalSpaces, Sikorski's Differentiable Spaces, Frolicher Spaces

## Definition 4.4: Let $X$ be a set.

A plot of dimension $p$ (or $p$-plot) on $X$ is a map from an open subset $O$ of $\mathbb{R}^{p}$ to $X$.

A diffeology on $X$ is a set $P$ of plots on $X$ such that, for all $\mathrm{pN},-$ any constant map $\mathbb{R}^{p} X$ is in $P$;

Let $I$ be an arbitrary set; let $\left\{f_{i}: O_{i} \rightarrow X\right\}_{i \in I}$ be a family of maps that extend to a map $f: \bigcup_{i \in I} O_{i} \rightarrow X$. If $\left\{f_{i}: O_{i} \rightarrow X\right\}_{i \in I} \subset P$, then $f \in P$. - (chain rule) Let $f \in P$, defined on $\in \mathbb{R}$. Let $q \in \mathbb{N}, O^{\prime}$ an open subset of Rq and $g$ a smooth map (in the usual sense) from $O^{\prime}$ to $O$. Then, $\mathrm{f} \circ \mathrm{g} \in \mathrm{P}$.

If $P$ is a diffeology $X,(X ; P)$ is called diffeological space. Let $(X ; P)$ et $\left(X^{\prime} ; P^{\prime}\right)$ be two diffeological spaces, a map $f: X \rightarrow X^{\prime}$ is differentiable (=smooth) if and only if $f(P) \subset P$.

Remark: Notice that any diffeological space $(X, P)$ can be endowed with the weaker topology such that all the maps that are in $P$ are continuous. But we prefer to mention this only for memory as well as other questions that are not closely related to our construction, and stay closer to the goals of this paper. Let us now define the Sikorski's differential spaces. Let $X$ be a Haussdorf topological space.

Definition 4.5: A (Sikorski's) differential space is a pair $(X ; F)$ where $F$ is a family of functions $X \rightarrow \mathbb{R}$ such that

- the topology of $X$ is the initial topology with respect to $F$
- for any $n \in N$, for any smooth function $\phi: \mathbb{R}^{n} \rightarrow \mathbb{R}$, for any $\left(f_{1}, \ldots, f_{n}\right) \in F^{\mathrm{n}}, \varphi \circ\left(f_{1}, \ldots, f_{n}\right) \in F$.

Let $(X ; F)$ et $\left(X^{\prime} ; F^{\prime}\right)$ be two differential spaces, a map $f: X \rightarrow X^{\prime}$ is differentiable (=smooth) if and only if $\mathrm{F}^{\prime} \circ f \in F$ :

We now introduce Frölicher spaces.
Definition 4.6: A Frölicher space is a triple $(X ; F ; C)$ such that - $C$ is a set of paths $\mathbb{R} \rightarrow X$,

- A function $f: X \rightarrow \mathbb{R}$ is in $F$ if and only if for any c $\in C$, $f \circ c \in C^{\infty}(\mathbb{R}, \mathbb{R})$;
- A path c : $\mathbb{R} \rightarrow X$ is in $C$ (i.e. is a contour) if and only if for any $f \in F, f \circ c \in C^{\infty}(\mathbb{R}, \mathbb{R})$.

Let $(X ; F ; C)$ et $\left(X^{\prime} ; F^{\prime} ; C^{\prime}\right)$ be two Frölicher spaces, a map $f: X \rightarrow X^{\prime}$ is differentiable (=smooth) if and only if $F^{\prime} \circ f \circ c \in C^{\infty}(\mathbb{R}, \mathbb{R})$

Any family of maps $F_{\mathrm{g}}$ from $X$ to $\mathbb{R}$ generate a Frolicher structure ( $X ; F ; C$ ), setting [14]:
$-C=\left\{c: \mathbb{R} \rightarrow X\right.$ Such that $\left.F_{g} \circ c \subset C^{\infty}(\mathbb{R}, \mathbb{R})\right\}$
$-F=\left\{f: \mathrm{X} \rightarrow \mathbb{R}\right.$ Such that $\left.f \circ c \subset C^{\infty}(\mathbb{R}, \mathbb{R})\right\}$
One easily see that $F_{\mathrm{g}} \subset F$. This notion will be useful in the sequel to describe in a simple way a Frolicher structure.

A Frölicher space, as a differential space, carries a natural topology, which is the pull-back topology of $\mathbb{R}$ via $F$. In the case of a finite dimensional differentiable manifold, the underlying topology of the Frölicher structure is the same as the manifold topology. In the infinite dimensional case, these two topologies differ very often.

In the three previous settings, we call $X$ a differentiable space, omitting the structure considered. Notice that, in the three previous settings, the sets of differentiable maps between two differentiable spaces satisfy the chain rule. Let us now compare these three settings: One can see (see e.g. [13]) that we have the following, given at each step by forgetful functors:
smooth manifold $\Rightarrow$ Frölicher space $\Rightarrow$ Sikorski differential space
Moreover, one remarks easily from the definitions that, if $f$ is a map from a Frölicher space $X$ to a Frölicher space $X^{\prime}$, f is smooth in the sense of Frölicher if and only if it is smooth in the sense of Sikorski.

One can remark, if $X$ is a Frölicher space, we define a natural diffeology on $X$ by Magnot [15]:
$P(F)=\coprod_{p \in \mathbb{N}}\left\{f p\right.$-paramatrization on $X ; F \circ f \in C^{\infty}(o, \mathbb{R})$ (in the usual sense) $\}$.

With this construction, we get also a natural diffeology when $X$ is a Frölicher space. In this case, one can easily show the following:

Proposition 4.7: Let $(X ; F ; C)$ and $\left(X^{\prime} ; F^{\prime} ; C^{\prime}\right)$ be two Frölicher spaces. A map $f: X \rightarrow X^{\prime}$ is smooth in the sense of Frölicher if and only if it is smooth for the underlying diffeologies [15].

Thus, we can also state:
smooth manifold $\Rightarrow$ Frölicher space $\Rightarrow$ Diffeological space

## Cartesian Products

The category of Sikorski differential spaces is not cartesianly closed, see e.g. [13]. This is why we prefer to avoid the questions related to cartesian products on differential spaces in this text, and focus on Frölicher and diffeological spaces, since the cartesian product is a tool essential for the definition of configuration spaces.

In the case of diffeological spaces, we have the following [12,16-19]:
Proposition 4.8: Let $(X ; P)$ and $\left(X^{\prime} ; P^{\prime}\right)$ be two diffeological spaces. We call product diffeology on $X \mathrm{x} X^{\prime}$ the diffeology $P \times P^{\prime}$ made of plots $g: O \rightarrow X \mathrm{x} X^{\prime}$ that decompose as $g=f \mathrm{x} f^{\prime}$, where $f: O \rightarrow X \in P$ and $f^{\prime}$ : $O \rightarrow X^{\prime} \in P^{\prime}$.

Then, in the case of a Frölicher space, we derive very easily, compare with e.g. Kriegl A et al. [14]:

Proposition 4.9: Let $(X ; F ; C)$ and $\left(X^{\prime} ; F^{\prime} ; C^{\prime}\right)$ be two Frölicher spaces, with natural diffeologies $P$ and $P^{\prime}$. There is a natural structure of Frölicher space on $X \mathrm{x} X^{\prime}$ which contours $C \mathrm{x} C^{\prime}$ are the 1-plots of $P \mathrm{x} P^{\prime}$.

We can even state the following results in the case of infinite products.

Proposition 4.10: Let I be an infinite set of indexes, that can be uncoutable.
(adapted from [21] ) Let $\{(X i ; P i)\}_{\text {iI }}$ be a family of diffeological spaces indexed by $I$. We call product diffeology on $\prod_{i \in I} X_{i}$ the diffeology $\prod_{i \in I} P_{i}$ made of plots $g: O \rightarrow \prod_{i \in I} X_{i}$ that decompose as $g=\prod_{i \in I} f_{i}$
where $f_{i} \in P_{i}$. This is the biggest diffeology for which the natural projections are smooth.

Let $\left\{\left(X_{i}, F_{i}, C_{i}\right)\right\}_{i} \in$ be a family of Frolicher spaces indexed by $I$, with natural diffeologies $P_{i}$. There is a natural structure of Frölicher space $\left(\prod_{i \in I} X_{i}, \prod_{i \in I} F_{i}, \prod_{i \in I} X C_{i}\right)$
which contours $\prod_{i \in I} X C_{i}$ are the 1-plots of $\prod_{i \in I} P_{i}$. A generating set of functions for this Frölicher space is the set of maps of the type:
$\varphi \circ \prod_{j \in J} f_{j}$
where $J$ is a finite subset of $I$ and $\varphi$ is a linear map $\mathbb{R}^{|J|} \rightarrow \mathbb{R}$.
Proof: By definition, following [12,20], $\prod_{i \in I} P_{i}$ is the biggest diffeology for which natural projections are smooth. Let $g: O \rightarrow X_{i}$ be a plot.

$$
g \in \varphi \circ \prod_{i \in I} p_{i} \Leftrightarrow p_{i} \circ g \in p_{i}
$$

where $p_{i}$ is the natural projection onto $X_{i}$, which gets the result.
With the previous point and Proposition 4.7, we get the family of contours of the product Frölicher space.

## Push-Forward, Quotient and Trace

We give here only the results that will be used in the sequel.
Proposition 4.11: Let $(X, P)$ be a diffeological space, and let $\mathrm{X}^{\prime}$ be a set. Let $f: X \rightarrow \mathrm{X}^{\prime}$ be a surjective map. Then, the set $[12,21]$.
$f(P)=\{u$ such that $u$ restricts to some maps of the type $f$ o $p ; p \in P\}$ is a diffeology on $X^{\prime}$, called the push-forward diffeology on $X^{\prime}$ by $f$.

We have now the tools needed to describe the diffeology on a quotient:

Proposition 4.12: Let $(X, P) \mathrm{b}$ a diffeological space and $R$ an equivalence relation on $X$. Then, there is a natural diffeology on $X / R$, noted by $P / R$, defined as the push-forward diffeology on $X / R$ by the quotient projection $X \rightarrow X / R$.

Given a subset $X_{0} \subset X$, where $X$ is a Frolicher space or a diffeological space, we can define on trace structure on $X_{0}$, induced by $X$.

If X is equipped with a diffeology $P$, we can define a diffeology $P_{\text {o }}$ on $X_{0}$ setting
$P_{0}=\left\{p \in P\right.$ such that the image of $p$ is a subset of $\left.X_{0}\right\}$.
If $(X, F, C)$ is a Frolicher space, we take as a generating set of maps $F_{g}$ on $X_{0}$ the restrictions of the maps $f \in F$. In that case, the contours (resp. the induced diffeology) on $X_{0}$ are the contours (resp. the plots) on $X$ which image is a subset of $X_{0}$.

## References

1. Fadell ER, Husseini SY (2012) Geometry and topology of configuration spaces. Springer Science \& Business Media, USA.
2. Ismagilov RS (1996) Representations of infinite-dimensional groups American Mathematical Soc.
3. Finkelstein D, Kondratiev Y, Kutoviy O (2010) Vlasov scaling for stochastic dynamics of continuous systems. J Stat Phys 141: 158-78.
4. Finkelstein D, Kondratiev Y, Kutoviy O (2012) Semigroup approach to birth-and-death stochastic dynamics in continuum. J Funct Anal 262: 1274-1308.
5. Finkelstein D, Kondratiev Y, Kozitsky Y, Kutoviy O (2011) Markov evolution of continuum particle systems with dispersion and competition.
6. Finkelstein D, Kondratiev Y, Kutoviy O (2013) Establishment and fecundity in spatial ecological models: statistical approach and kinetic equations. Infinite Dimensional Analysis, Quantum Probability and Related Topics 16: 1350014.
7. Albeverio S, Daletskii A, Lytvynov E (2001) De Rham cohomology of configuration spaces with Poisson measure. J Funct Anal 185: 240-273.
8. Hilborn $R$ (2012) Chaos an non linear dynamics, Oxford university Press UK
9. Hagedorn D, Kondratiev Y, Pasurek T, Röckner M (2013) Gibbs states over the cone of discrete measures. J Funct Anal 264: 2550-2583.
10. Hughes T, Marsden JE (1983) Mathematical foundations of elasticity PrenticeHall Civil Engineering and Engineering Mechanics Series. Englewood Clis, New Jersey, Prentice-Hall, USA.
11. Ohta SI (2014) Examples of spaces with branching geodesics satisfying the curvature-dimension condition. Bull Lond Math Soc 46: 19-25.
12. Souriau JM (1986) Un algorithme générateur de structures quantiques.
13. Cherenack P, Ntumba P (2001) Spaces with differential structure and an application to cosmology. Demonstr Math 34: 161-180.
14. Kriegl A, Michor PW (1997) The convenient setting of global analysis. American Mathematical Soc, USA.
15. Magnot JP (2008) Difféologie du fibré d'Holonomie en dimension infinite. Math Rep Can Roy Math Soc 28: 121-128.
16. Donato P (1984) Revetements de groupes differentiels These de doctorat detat. Universite de Provence, Marseille, UK.
17. Frölicher A, Kriegl A (1988) Linear spaces and differentiation theory, John Wiley \& Sons Inc, USA.
18. Iglesias $P$ (1987) Connexions et diffeologie Aspects dynamiques et topologiques des groupes infinis de transformation de la mecanique Travaux en cours 25: 61-78.
19. Leslie J (2003) On a diffeological group realization of certain generalized symmetrizable Kac-Moody Lie algebras. J Lie Theor 13: 427-442.
20. Lesne A (1996) Méthodes de renormalisation. Phénomènes critiques, chaos, structures fractales, Eyrolles, Paris, France.
21. Magnot JP (2013) Ambrose-Singer Theorem on Diffeological Bundles and Complete Integrability of the Kp Equation. Int J Geom Methods Mod Phys 10: 1350043.

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[^0]:    *Corresponding author: Jean-Pierre Magnot, Department of Mathematics, University of Angers, Lycee Jeanne dArc, 30 avenue de Grande Bretagne, F-63000, Clermont-Ferrand, France, Tel: +33 2419623 23; E-mail: jean-pierr.magnot@ac-clermont.fr
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